

Universal, fault-tolerant gate sets

Recall:

and efficient 😊

① The Clifford group "is" FT but *classically simulable*

② Transversal gates can't get us to universality

How can we proceed?

We proceed by *relaxing* our definition of universality, and then hoping we can figure out how to scrape by with a non-transversal gate...

Explicitly, we will show that

$$C_n \cup \{T\}$$

suffices with this relaxed notion, and then give fault-tolerant constructions of the T gate

Approximating unitaries

Given a **finite** set of gates G , $\langle G \rangle$ is **countable**, but $U(n)$ is **uncountable**
 $\uparrow$ **$n \times n$ unitary matrices**
 $\swarrow$ **circuits over G**

Since we can't hope to implement every unitary exactly, we need to have a notion of **approximation error**.

(Approximation error)

Given a unitary U , the **error** when U is approximated by V is

$$\|U - V\| = \max_{|\psi\rangle} \| (U - V)|\psi\rangle \|$$

$\nwarrow$ **operator norm**

Note: $\|\cdot\|$ is a norm, so $\|A + B\| \leq \|A\| + \|B\|$
also $\|UA\| = \|A\|$, U unitary

We use the operator norm because the difference in measurement probabilities after apply U or V is bounded by $2\|U - V\|$

Def'n

A finite set of gates G , G is **universal** for a quantum computation if for any unitary U and error rate $\epsilon > 0$, there exists a circuit V over G such that

$$\|U - V\| \leq \epsilon$$

Approximation of single-qubit unitaries suffices

We first show that it suffices to only approximate **single-qubit unitaries** assuming CNOT is given. For this it is helpful to note an important fact about how errors in a circuit combine

(Approximation errors add) — important!

Observe that if U_1 and U_2 are approximated to error ϵ by V_1 and V_2 , resp., then

$$\begin{aligned}\|U_1 U_2 - V_1 V_2\| &= \|U_1 U_2 - U_1 V_2 + U_1 V_2 - V_1 V_2\| \\ &= \|U_1 (U_2 - V_2) + (U_1 - V_1) V_2\| \\ &\leq \|U_2 - V_2\| + \|U_1 - V_1\| \\ &= 2\epsilon\end{aligned}$$

The preceding point is useful in compilation because it tells us that if we want to approximate a circuit with d gates to error ϵ , we need to approximate each gate to error ϵ/d

Prop.

Given a set of single-qubit gates G which is universal for 1-qubit unitaries, $G \cup \{\text{CNOT}\}$ is universal for n -qubit unitaries.

Pf.

Write a unitary over CNOT + single qubit rotations, then approximate each to error ϵ/d

Universal sets of quantum gates

Thm.

$\{H, T\}$ is universal for single-qubit unitaries

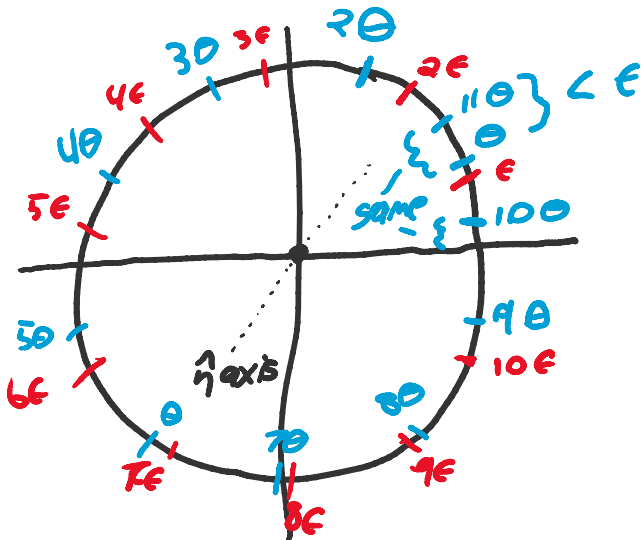
Pf. Sketch

Let $R_1 = THTH$, $R_2 = HR_1H = HTHT$.

It can be shown that R_1 & R_2 are rotations of an angle θ which is an irrational multiple of 2π (i.e. $k\theta \neq 2\pi \forall k \in \mathbb{Z}$)

around non-parallel axes $\hat{n}_1, \hat{n}_2$ resp.

Since θ is irrational, we can rotate by infinitely many angles around either axis. Intuitively, say we want to approximate $R_{\hat{n}}(\alpha)$ to accuracy ϵ .



$$\therefore \|R_{\hat{n}}(100\theta) - I\| < \epsilon$$

so multiples of $R_{\hat{n}}(100\theta)$
fill the space with points
 $< \epsilon$ distance apart
(called an ϵ -net)

So write any single-qubit unitary as

$$R_{\hat{n}}(\beta) R_{\hat{m}}(\delta) R_{\hat{n}}(\gamma)$$

and approximate each 😊

Speed of approximation

Naively, this method generates an ϵ -net with circuits of length $O(1/\epsilon^2)$ - roughly $2\pi/\epsilon$ rotations to get within ϵ of I , then $2\pi/\epsilon$ multiples to cover the circle.

This sucks!

Thm. (Solovay-Kitaev) *proven unnecessary just last year!!!*

Given a gate set G closed under inverses that is universal for single-qubit unitaries, a single-qubit unitary U can be approximated over G to ϵ error with at most $O(\log^c(1/\epsilon))$ gates from G

The Solovay-Kitaev theorem is fundamental to the possibility of FTQEC, as it states that we can use fault-tolerant gates with sub-polynomial overhead. In practice it is important for FTQEC implementations of

- The Quantum Fourier transform (basis of Shor + others)
- Hamiltonian simulation

both of which require many single-qubit gate approximations.

Other universal gate sets

Thm.

The Clifford group plus any non-Clifford unitary is approximately universal

(The proof of this theorem brings up interesting connections to lattice theory. Certain gate sets correspond to automorphisms of well-known lattices. In the Clifford case, this is the Barnes-Wall lattice, and the most commonly cited proof for this fact falls out of this connection)

The following are ^{some} approximately universal gate sets and their inclusions:

$$\{H, CNOT, T\}$$

⊂

⊂

⊂

roughly

$$\{H, CH, Toffoli\}$$

$$\{e^{i\pi/4}H, S, Toffoli\}$$

$$\{\sqrt{H}, Toffoli\}$$

also this

⊂

⊂

⊂

These mean that complex numbers are NOT necessary for QC!

$$\{H, Toffoli\} = \{H, CCZ\}$$

1U (becomes equality with ancillas)

$$\{H \otimes H, X, CX, CCX\}$$

↑

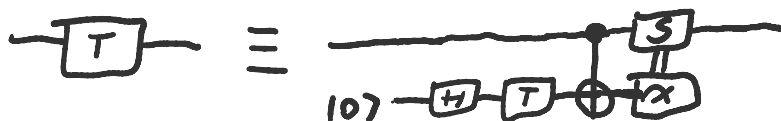
In fact, not even Real numbers are needed! QC can be done with finite precision!

Non-transversal T gates

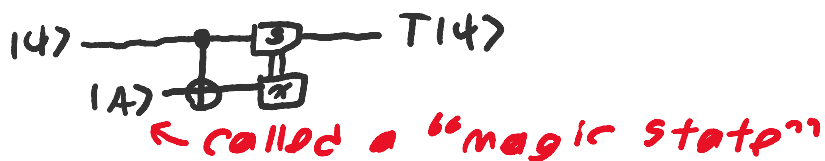
We know we need at least one gate (T) which is **NOT** transversal to get a universal gate set. Fortunately there exist other **more costly** methods for **certain gates**, namely **Gate teleportation + magic states**.

(Teleporting the T gate)

The following circuits are equivalent up to global phase



If we write $|A\rangle = TH|0\rangle$, then up to global phase



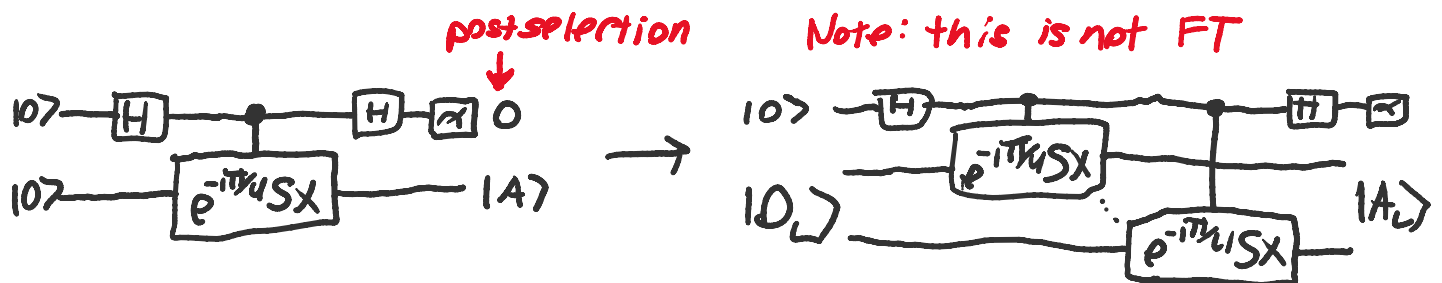
Upshot: if we can prepare the state $|A\rangle$, perform Clifford gates, and measure fault-tolerantly, we can perform the T gate fault-tolerantly!

($|A\rangle$ states circa early 2000's)

The $|A\rangle$ state can be prepared fault-tolerantly by using transversal Cliffords. In particular,

$$e^{-i\pi/4} SX |A\rangle = |A\rangle$$

So we can **project** a state onto the $|A\rangle$ state by measuring $e^{-i\pi/4} SX$, which is Clifford. Explicitly,



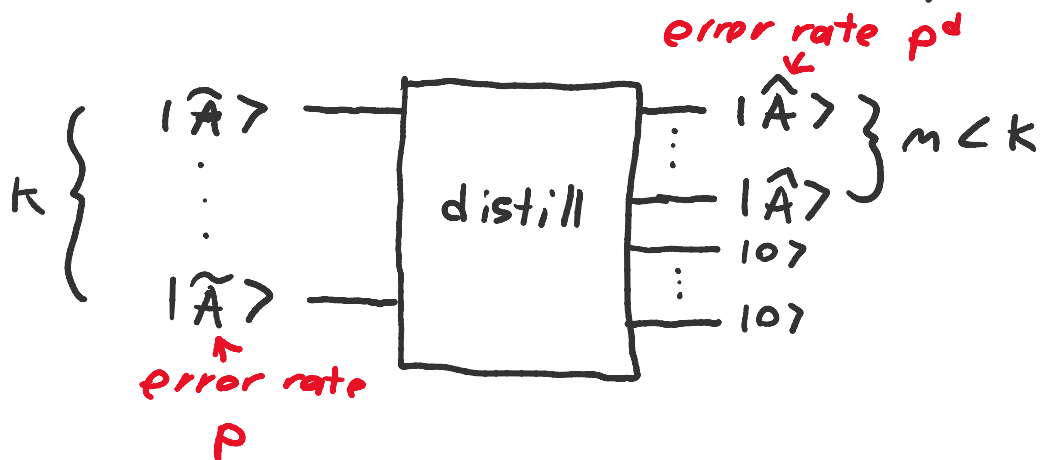
State distillation

In many modern codes, the old-style $|A\rangle$ state preparation is very inefficient.

Bravyi & Kitaev (2005) gave an alternative method which scales better called **magic state distillation**

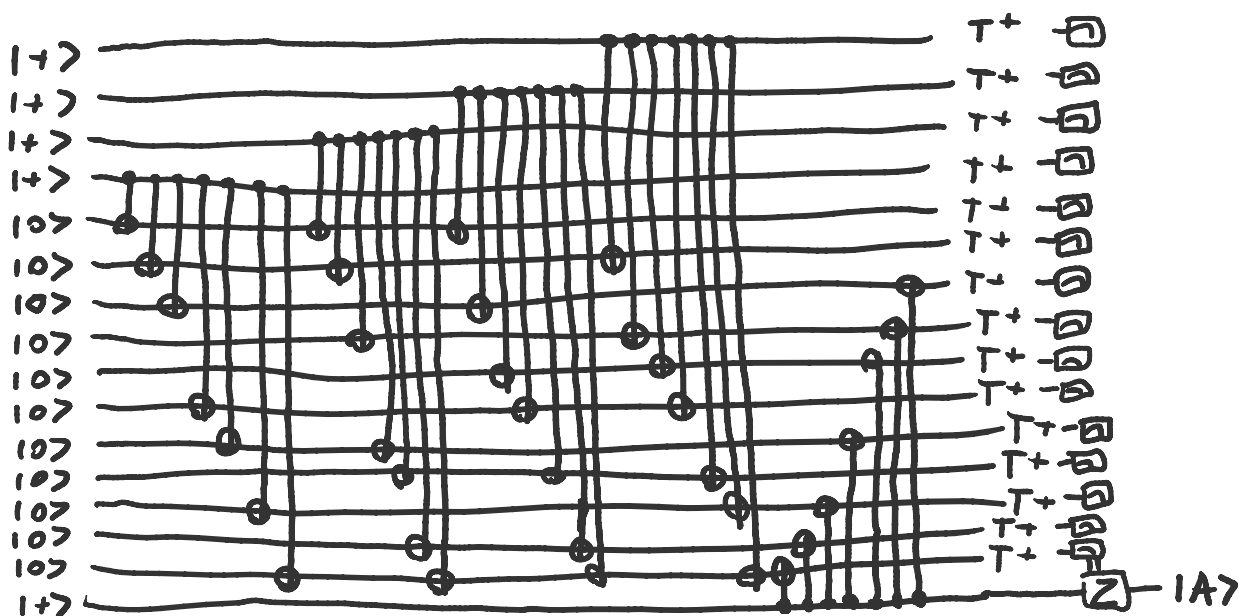
(magic state distillation)

The basic idea of magic state distillation is to take many **faulty** copies of the magic state, and distill them into **one or a few** better copies



(15-1 protocol)

The 15-1 protocol uses a (punctured) Reed-Muller code for which the T gate is transversal to get **cubic** error reduction (i.e. $p \xrightarrow{\text{distill}} O(p^3)$). The details are unimportant, but the **time and space** it uses is!

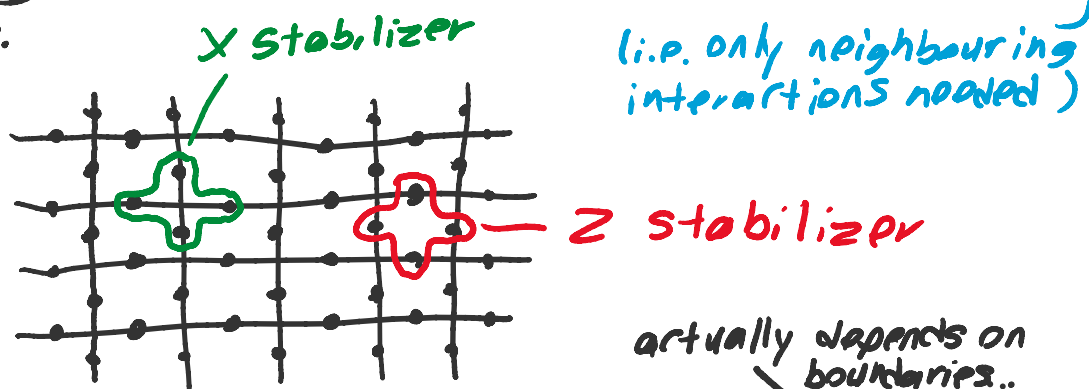


Topological codes

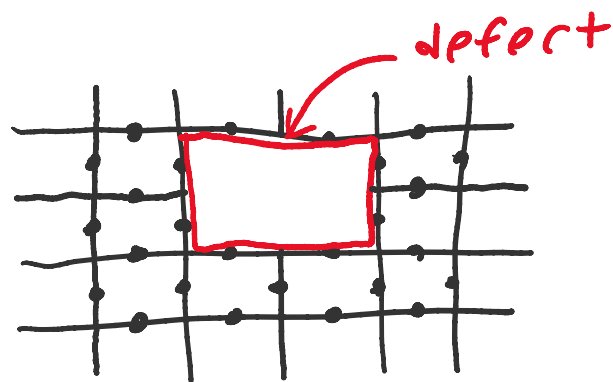
Since the 2010's, Kitaev's **Surface** and more generally **topological codes** have, with the help of state distillation, supplanted the leading CSS code (**Steane code**).

(Surface code)

The surface code is a stabilizer code laid out on a 2-D lattice. The generators are importantly **local**, corresponding to X 's around each vertex and Z 's around each face.



On a square lattice, this gives **no logical qubits** so logical qubits are made by cutting **defects** from the stabilizer generators



Defects can then be **moved** with only measurement of the stabilizers

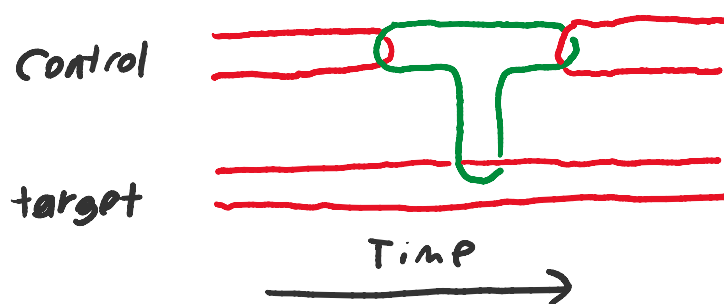
Fault-tolerant topological gates

Fault-tolerant gates in a topological QECC have different flavours, but generally **Clifford gates are CHEAP**, **T gates are EXPENSIVE**. In particular, Clifford gates are generally performed by moving defects, while T gates use magic state distillation.

(Braided surface code - until 2020's)

The first main FTQEC proposal based on surface codes. Defects (**primal** or **dual**) are in pairs, implement gates by braiding the defects

e.g. CNOT:



"only the topology matters"



time could be vertical...

spawned compilation work looking at topological computation

(Lattice surgery - the new kid)

Leading proposals now use lattice surgery.

Main idea is to compute (Cliffords) by **merging** and **splitting** defects.



Relative Volumes

Computation in topological codes is typically measured in **volume**. By volume, state distillation vastly outweighs computation!

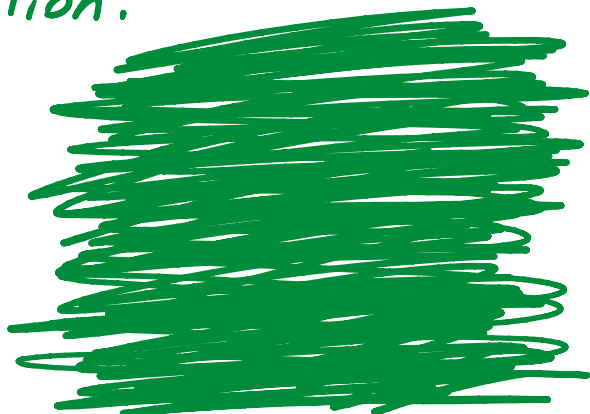
E.g.



X gate

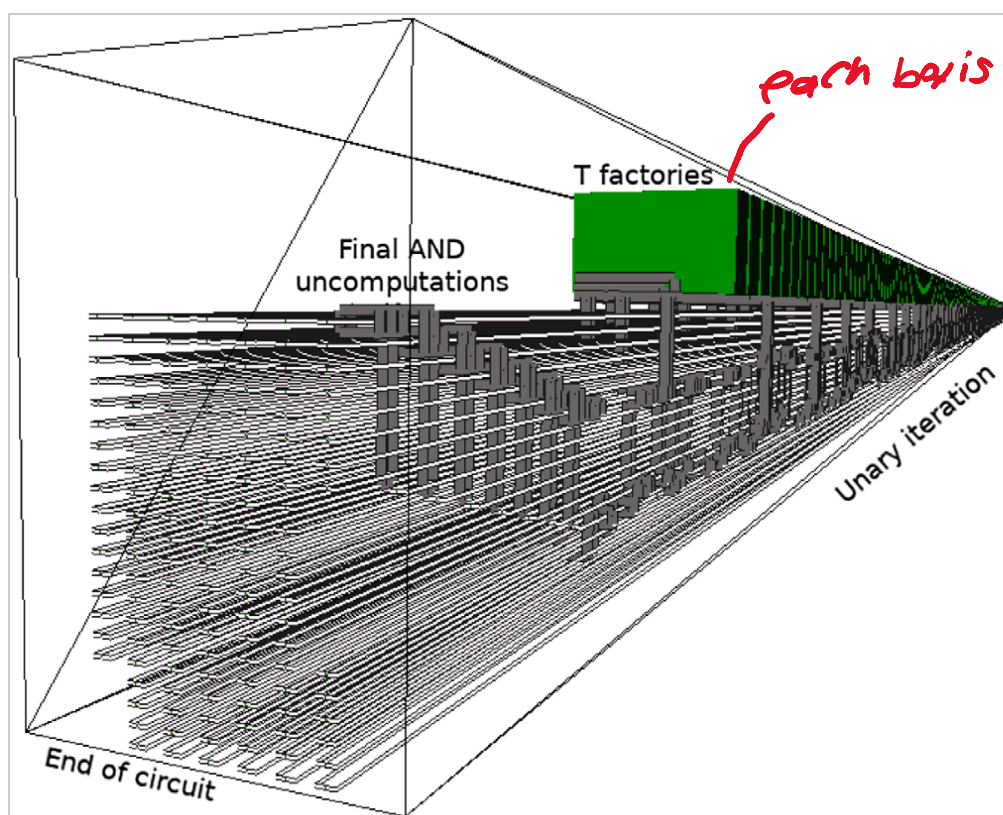
CNOT gate

$|A\rangle$ state



Result: The T gate is **ORDERS OF MAGNITUDE** more expensive in surface codes than Clifford gates

E.g. from arXiv:1805.03662, Hamiltonian simulation with **linear** T-complexity



Further notes about resources

In FTQEC, **space** and **time** usage are linked and increases in either at the logical layer tend to compound and affect both physically.

Thm. (Threshold theorem)

A threshold theorem for a FTQEC method states that if **physical error rates** are below p_{th} , then a computation can be performed to arbitrary precision by increasing the amount of error correction.

Ex.

p_{th} for the **Steane code** is proven to be 10^{-5}

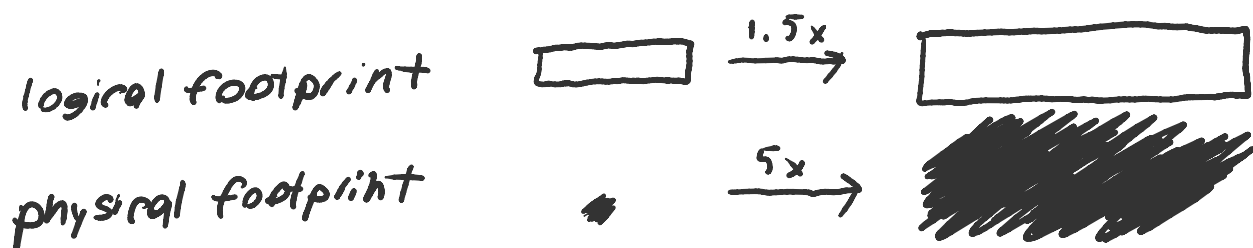
p_{th} for **surface code** is estimated at 10^{-2}

(Code distance)

The distance of a QEC is a measure of how resilient to errors it is. Increasing distance requires (typically) increasing physical qubits per logical qubit.

To understand how resource usage compounds in FTQEC, imagine that one circuit C_1 takes twice as long (i.e. depth) as another circuit C_2 . When applying error correction, C_1 has more chance for errors to accumulate, and so may take more physical space as well. If C_1 has twice as many T gates as well, each needs a magic state which is 50% more precise, leading to larger & longer state preparation circuits.

e.g. $15-1$ becomes $15^2-1...$



Beyond Solovay-Kitaev

Since gate approximation produces **long** circuits with **many T gates**, optimizing single qubit gate approximations in Clifford+T is a fundamental problem in circuit compilation

For many years, the best-known **Constructive** method (Dawson & Nielsen 2005) gave $C > 3$ for $\{H, T\}$. An information-theoretic lower bound of $3 \log(1/\epsilon)$ suggested better was possible but wasn't met until the **Number theoretic method** was developed in 2013-2014. To this day, for $\{H, T\}$ the best known (**pure circuit**) algorithm is Ross & Selinger (2014)'s **Grid-Synth** which produces approximations of length $3 \log(1/\epsilon) + O(\log \log 1/\epsilon)$

If measurement is allowed, the **PQF** method of Bocharov, Roetteler & Svore (2015) reduces this to an expected length

$$\log(1/\epsilon) + O(\log \log 1/\epsilon)$$

